

FREAK WAVES AND STOCHASTIC WAVE GROUPS



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A NATURAL BEAUTY !





Freak waves



Rogue waves



Giant waves



Extreme waves



Rogue waves



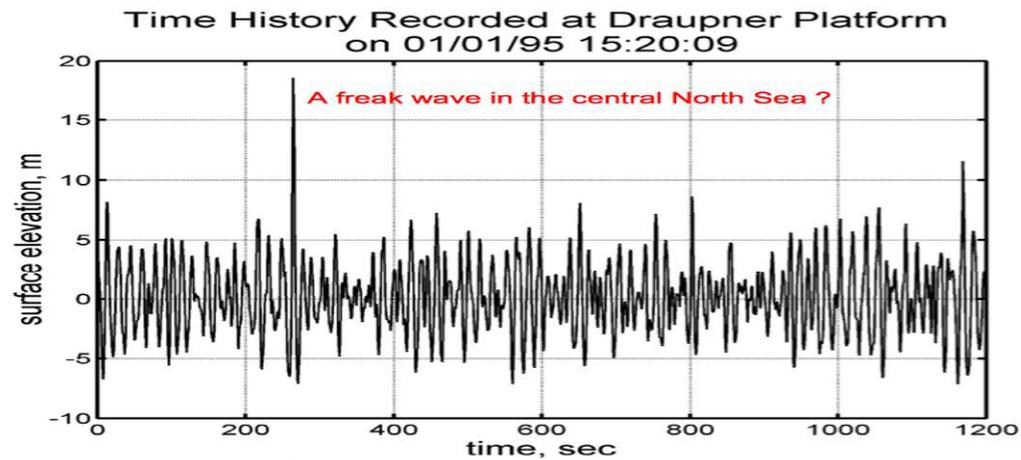
Extreme waves

Giant waves

Freak waves



DRAUPNER EVENT JANUARY 1995



$$H_{\max} = 25.6 \text{ m}$$

with

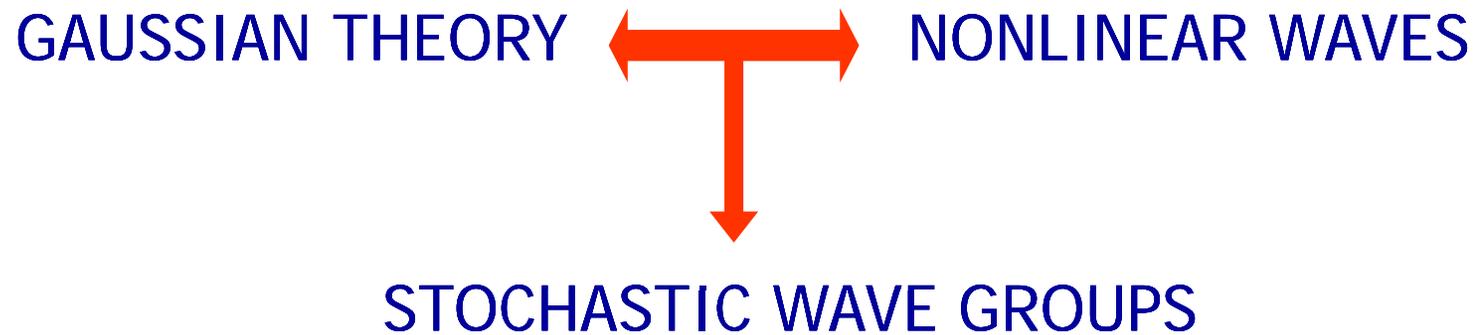
$$H_{\max}/H_s > 2 !$$

1 in 200,000 waves

according to Rayleigh law



CAN WE MODEL AND PREDICT EXTREME WAVES ?



STOKES EQUATIONS

$$\frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial z^2} = 0$$

$$v_z = \frac{\partial \Phi}{\partial z}$$

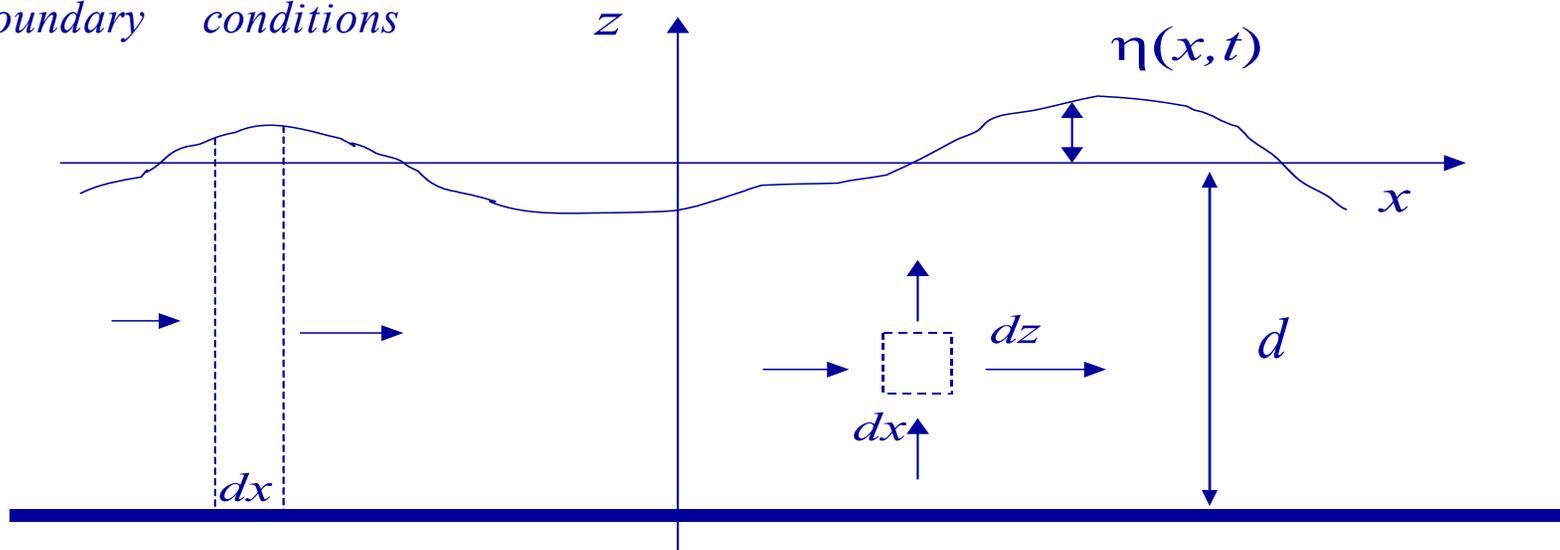
$$v_x = \frac{\partial \Phi}{\partial x}$$

$$\left(\frac{\partial \Phi}{\partial z} \right)_{z=\eta} = \frac{\partial \eta}{\partial t} + \frac{\partial \eta}{\partial x} \left(\frac{\partial \Phi}{\partial x} \right)_{z=\eta}$$

Inviscid, irrotational

$$\left(\frac{\partial \Phi}{\partial t} \right)_{z=\eta} + \frac{1}{2} \left[\left(\frac{\partial \Phi}{\partial x} \right)^2 + \left(\frac{\partial \Phi}{\partial z} \right)^2 \right]_{z=\eta} + g\eta = f(t)$$

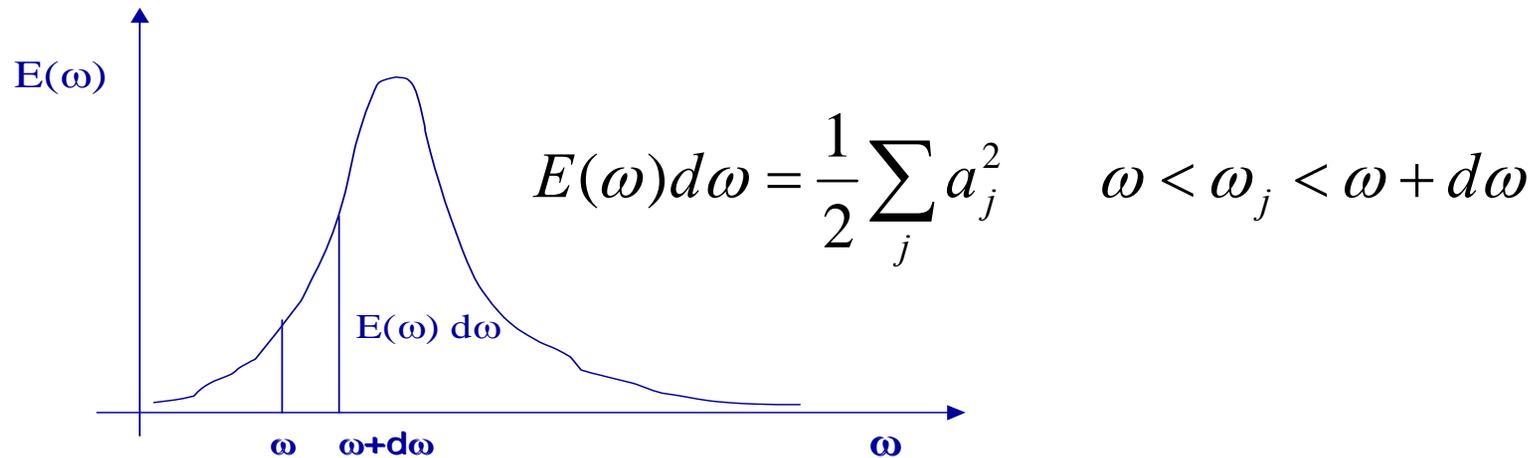
boundary conditions



LINEAR WAVES : GAUSSIAN SEAS

$$\eta(x, t) = \sum_{j=1}^N a_j \cos(k_j x + \omega_j t + \varepsilon_j)$$

ε_j uniformly random in $[0, 2\pi]$



$$\psi(T) = \langle \eta(x_0, t_0) \eta(x_0, t_0 + T) \rangle = \int_0^{\infty} E(\omega) \cos \omega T d\omega$$

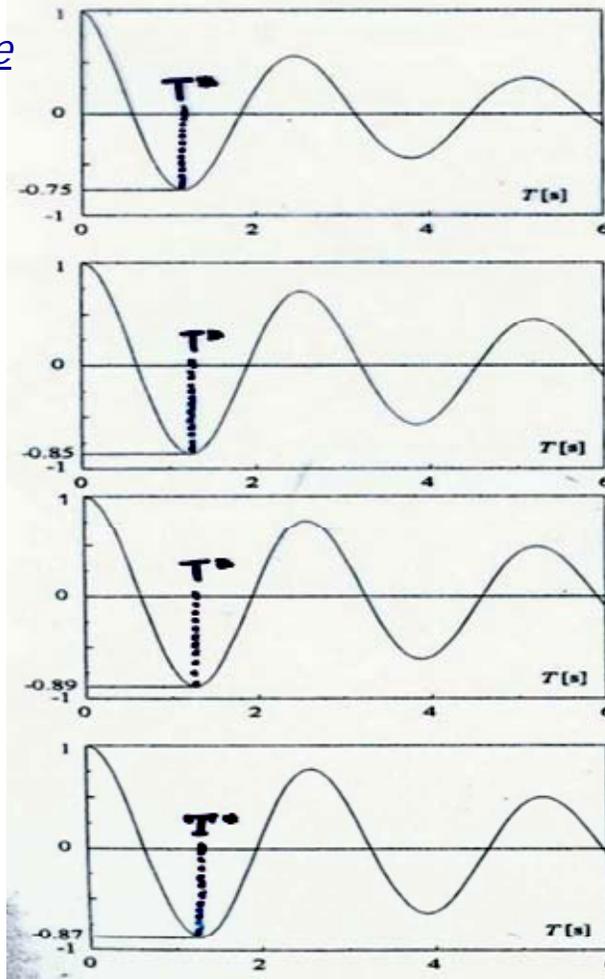
Stationarity

Ergodicity

Gaussianity

TYPICAL WAVE SPECTRA OF THE MEDITERRANEAN SEA*

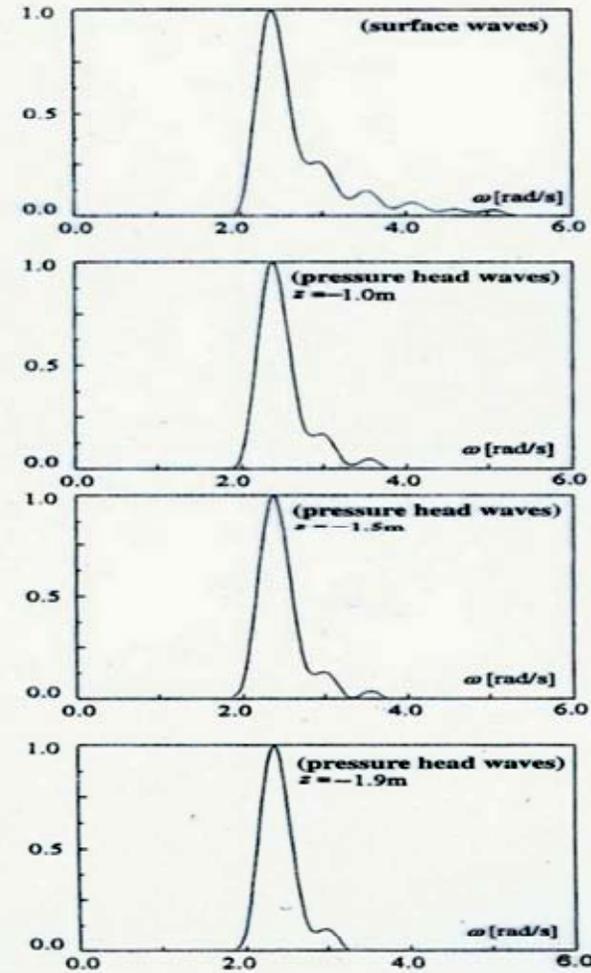
Time covariance



Wind generated waves: basic concepts

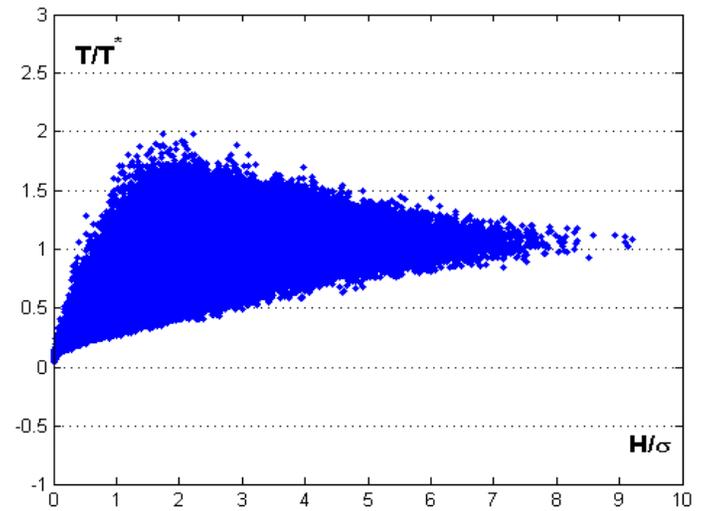
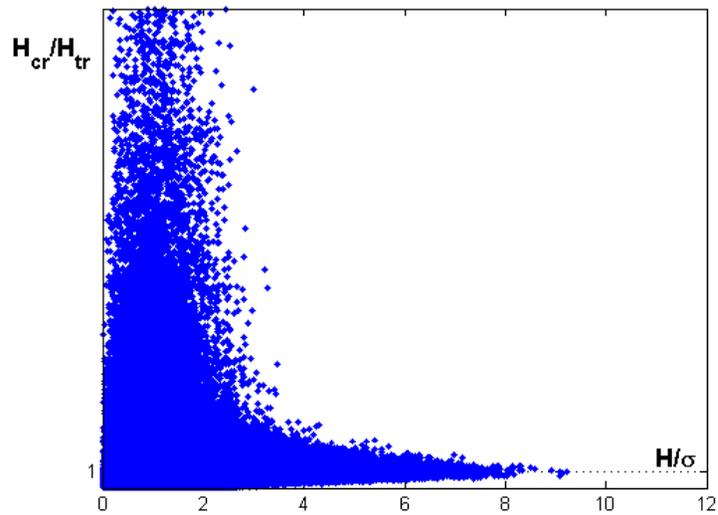
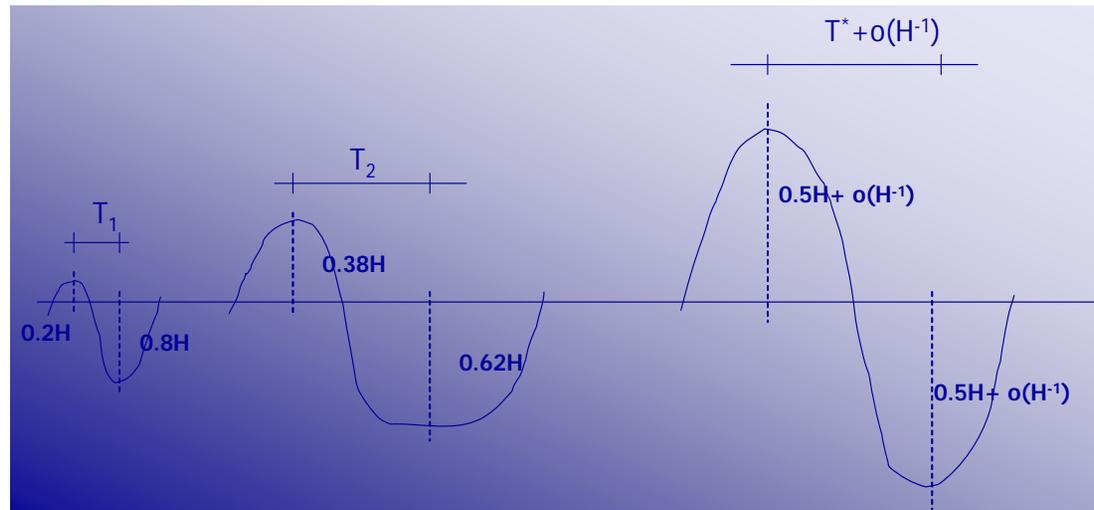
143

Spectrum



*from *Boccotti P. Wave Mechanics 2000 Elsevier*

OCCURRENCE OF A HIGH WAVE IN GAUSSIAN SEAS*



*Theory of quasi-determinism, *Boccotti P. Wave Mechanics 2000 Elsevier*

EXPECTED SHAPE OF THE SEA LOCALLY TO A HIGH WAVE*

What happens in the neighborhood of a point \mathbf{x}_0 if a large crest followed by large trough are recorded in time at \mathbf{x}_0 ?

$$\Pr \left[\begin{array}{l} \eta(\mathbf{x}_0 + \mathbf{X}, t_0 + T) \in (u, u + du) \\ \text{conditione } d \text{ to} \\ \eta(\mathbf{x}_0, t_0) = H/2, \eta(\mathbf{x}_0, t_0 + T_2^*) = -H/2 \end{array} \right]$$

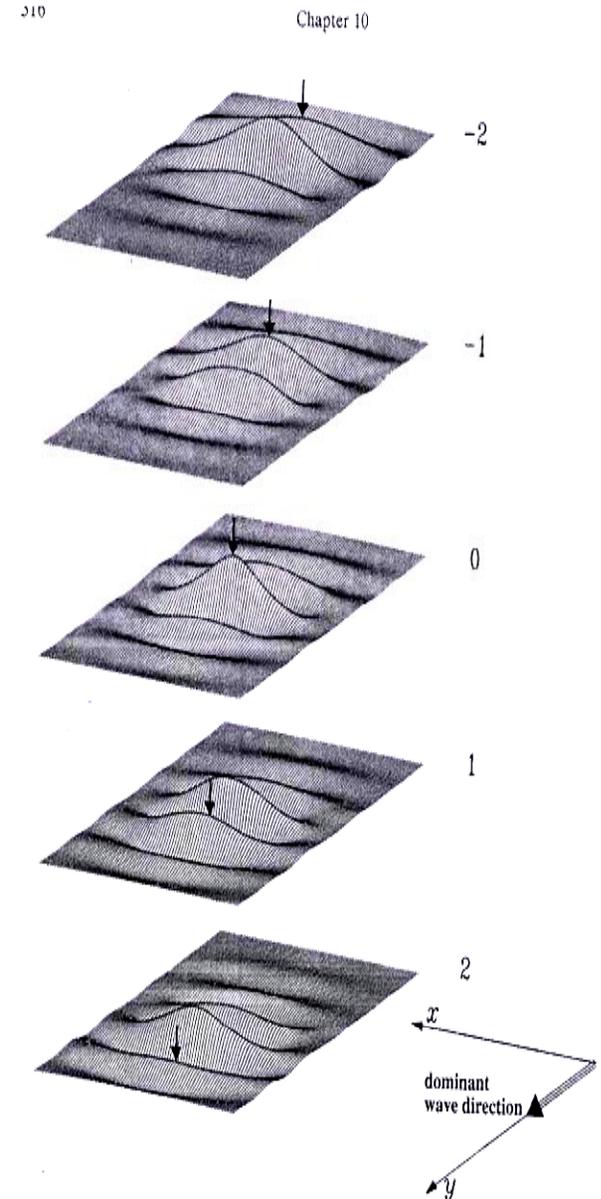
$$\Downarrow \frac{H}{\sigma} \rightarrow \infty$$

$$\eta_c(\mathbf{x}_0 + \mathbf{X}, t_0 + T) = \frac{H}{2} \Psi(\mathbf{X}, T)$$

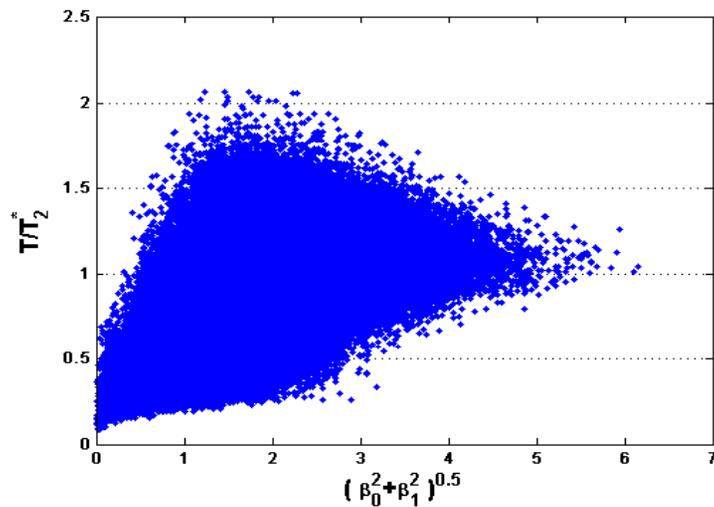
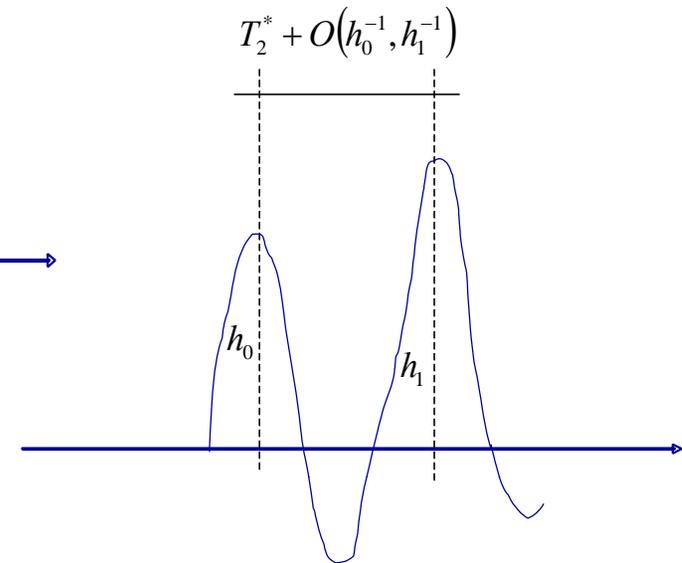
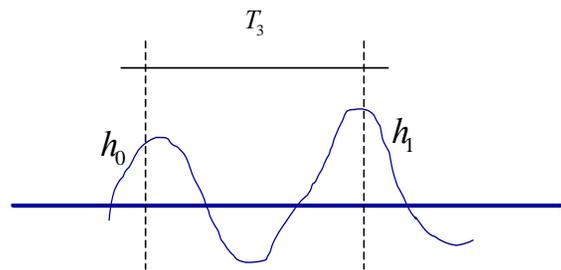
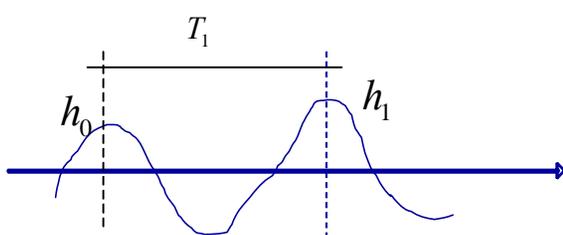
SPACE-TIME covariance

$$\Psi(\mathbf{X}, T) = \langle \eta(\mathbf{x}_0, t_0) \eta(\mathbf{x}_0 + \mathbf{X}, t_0 + T) \rangle$$

* Boccotti P. **Wave Mechanics** 2000 Elsevier



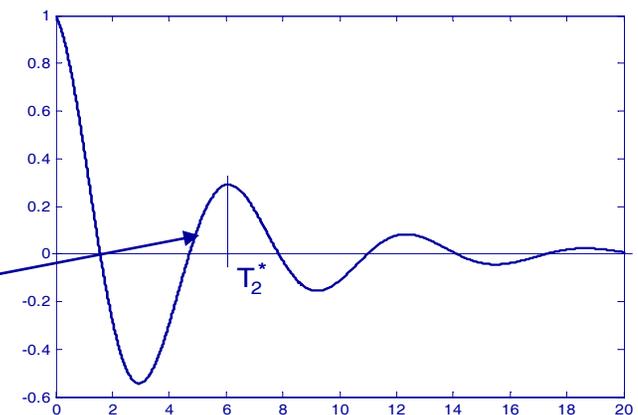
SUCCESSIVE WAVE CRESTS IN TIME*



as $h_0 \rightarrow \infty, h_1 \rightarrow \infty$

$$p(\beta_0, \beta_1) = \frac{\beta_0 \beta_1}{1 - k^2} \exp\left[-\frac{\beta_0^2 + \beta_1^2}{2(1 - k^2)}\right] I_0\left(\frac{k\beta_0\beta_1}{1 - k^2}\right)$$

Bivariate Rayleigh



* Fedele F., **Successive wave crests in a Gaussian sea**, *Probabilistic Eng. Mechanics* 2005 vol. 20, Issue 4, 355-363

EXPECTED SHAPE OF THE SEA LOCALLY TO TWO SUCCESSIVE WAVE CRESTS *

What happens in the neighborhood of a point \mathbf{x}_0
if two large successive wave crests are recorded in time at \mathbf{x}_0 ?

$$\Pr \left[\begin{array}{l} \eta(\mathbf{x}_0 + \mathbf{X}, t_0 + T) \in (u, u + du) \\ \text{conditioned to } \eta(\mathbf{x}_0, t_0) = h_1, \eta(\mathbf{x}_0, t_0 + T_2^*) = h_2 \end{array} \right]$$

$$\Downarrow \quad \frac{h_1}{\sigma} \rightarrow \infty, \quad \frac{h_2}{\sigma} \rightarrow \infty$$

$$\eta_c(\mathbf{X}, T) = \frac{\psi(0)h_1 - h_2\psi(T_2^*)}{\psi^2(0) - \psi(T_2^*)^2} \Psi(\mathbf{X}, T) + \frac{\psi(0)h_2 - h_1\psi(T_2^*)}{\psi^2(0) - \psi(T_2^*)^2} \Psi(\mathbf{X}, T - T_2^*)$$

What is hidden “behind” this equation ?

* Fedele F., 2006. **On wave groups in a Gaussian sea.** *Ocean Engineering 2006 (in press)*

A SINGLE WAVE GROUP CAUSES TWO SUCCESSIVE WAVE CRESTS !*

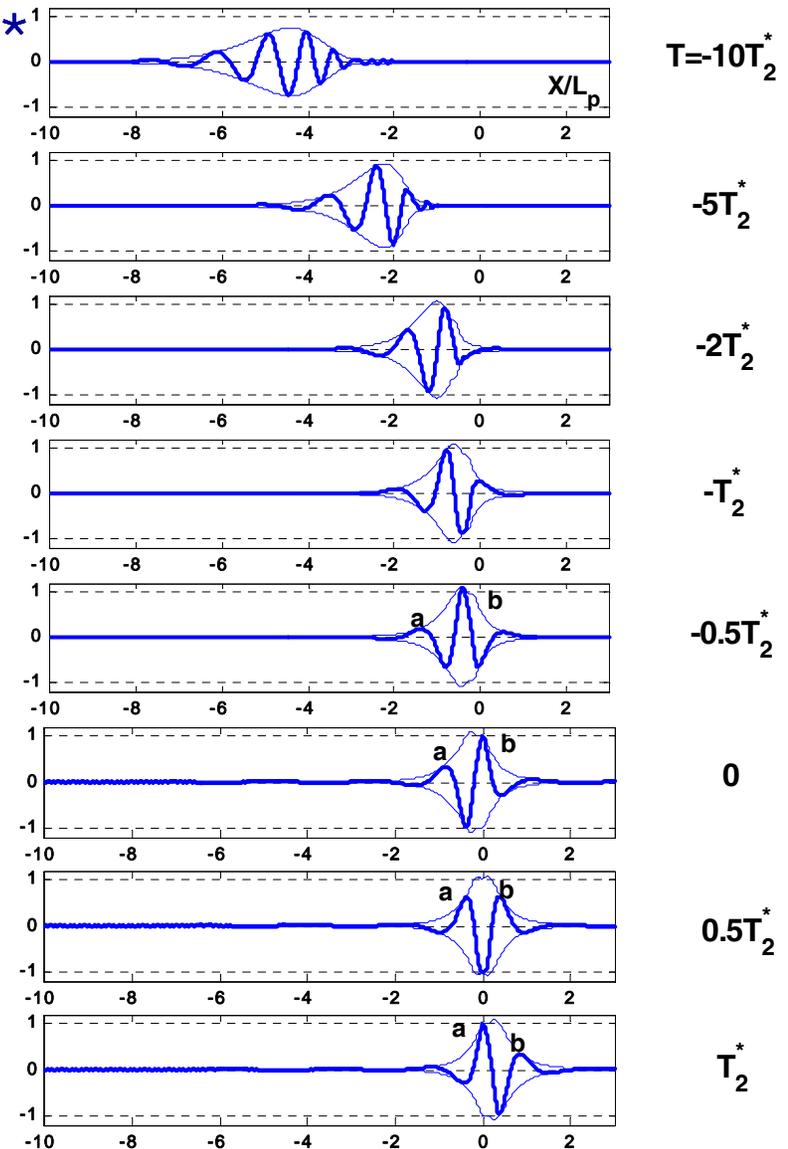
$$\eta_{ex}(X, T) = h \Psi(X, T)$$

SPACE-TIME covariance

STOCHASTIC WAVE GROUP

amplitude h random variable
distributed according to Rayleigh

stochastic family of wave groups

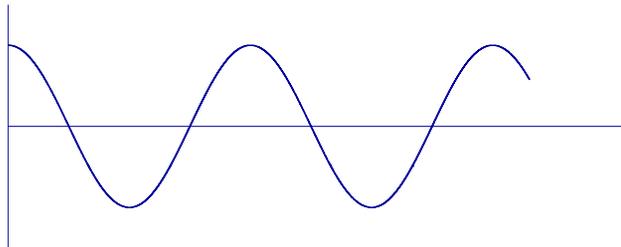


* Fedele F., 2006. **On wave groups in a Gaussian sea.** *Ocean Engineering 2006 (in press)*

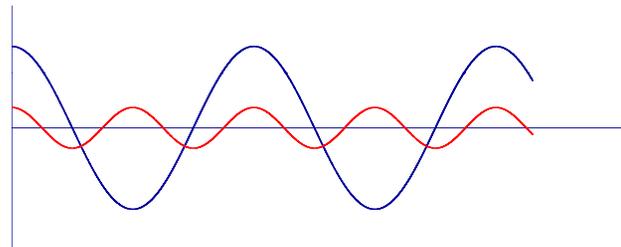
NONLINEAR RANDOM SEAS

Second order effects: **BOUND WAVES**

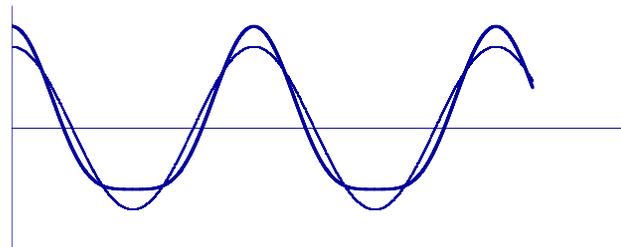
LINEAR TERM



LINEAR & NON-LINEAR TERMS



NON LINEAR WAVE

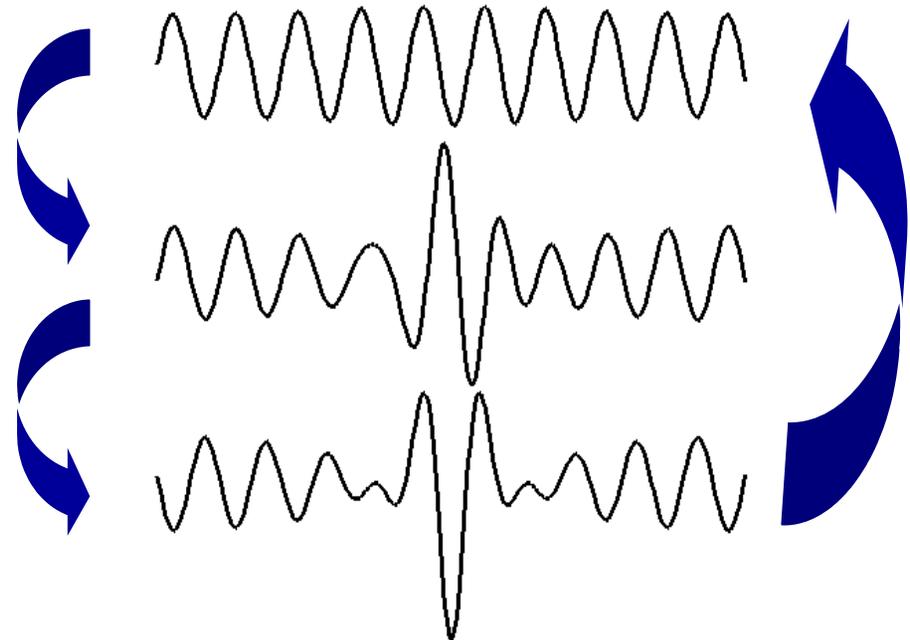


Crest-trough asymmetry : **skewness** >0

Wave height almost linear

Effects on Short time scale : wave period

Third order effects : **FOUR-WAVE RESONANCE**



Crest-trough symmetry : **kurtosis** >3

Benjamin-Feir Index $BFI = \text{steepness} / \text{bandwidth}$

Modulation instability: Fermi Ulam-Pasta recurrence

Effects on slow time scale : wave period/ steepness^2

**DOMINANT ONLY IN
UNIDIRECTIONAL NARROW-BAND SEAS !**

NONLINEAR RANDOM SEAS cont'd

$$\tilde{\eta}(\underline{\mathbf{x}}, t) = \eta(\underline{\mathbf{x}}, t) + \eta_2(\underline{\mathbf{x}}, t)$$

Third order effects :
FOUR-WAVE RESONANCE
 (WEAK WAVE TURBULENCE)

$$\eta(\underline{\mathbf{x}}, t) = \frac{1}{\pi} \sum_n \sqrt{\frac{\omega_n}{2g}} |B_n(t)| \cos(\underline{\mathbf{k}}_n \cdot \underline{\mathbf{x}} + |\varphi_n(t)|)$$

$$\frac{dB_n}{dt} + i\omega_n B_n = -i \sum_{p,q,r} T_{npqr} \delta_{n+p-q-r} B_p^* B_q B_r$$

Conserved quantities :

Hamiltonian

Wave action

Wave momentum

$$H = \sum_n \omega_n B_n(t) B_n^*(t) + \frac{1}{2} \sum_{n,p,q,r} T_{npqr} \delta_{n+p-q-r} B_n^*(t) B_p^*(t) B_q(t) B_r(t)$$

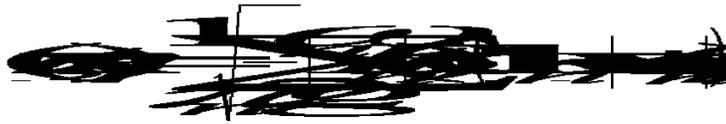
$$\mathbf{A} = \sum_n B_n(t) B_n^*(t) \quad \mathbf{M} = \sum_n \mathbf{k}_n B_n(t) B_n^*(t)$$

Second order effects:
BOUND WAVES

$$\eta_2(\underline{\mathbf{x}}, t) = \sum_{p,q} \sqrt{\omega_p \omega_q} |B_p(t)| |B_q(t)| \Gamma_{pq}^+ \cos((\underline{\mathbf{k}}_p + \underline{\mathbf{k}}_q) \cdot \underline{\mathbf{x}} + \varphi_p + \varphi_q) +$$

$$+ \sum_{p,q} \sqrt{\omega_p \omega_q} |B_p(t)| |B_q(t)| \Gamma_{pq}^- \cos((\underline{\mathbf{k}}_p - \underline{\mathbf{k}}_q) \cdot \underline{\mathbf{x}} + \varphi_p - \varphi_q)$$

NONLINEAR EVOLUTION OF A STOCHASTIC WAVE GROUP*



At $(x=0, t=0)$ we impose that all the elementary waves are in phase (focussing)

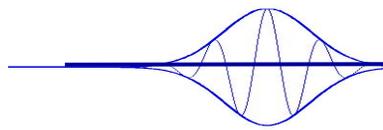
$$\varphi_n(0) = 0 \quad n = 1, \dots, N$$

Hamiltonian, wave action
and wave momentum
are conserved

t=0

t=-t0

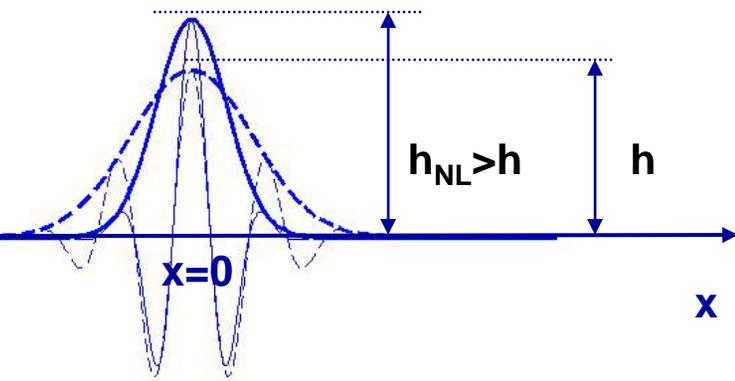
(linear wave group)



$$h_{NL} = (1 + \lambda)h$$

$$\tilde{h}_{nl} = (1 + \lambda)h + \alpha(1 + \lambda)^2 h^2$$

h Rayleigh distributed



λ Self-focussing parameter
 α Skewness parameter

$$\Pr(H_{\max} > h) = \exp \left[-\frac{(1 + \lambda)^2}{8\alpha^2} \left(1 - \sqrt{1 + \frac{4\alpha}{(1 + \lambda)^2} \frac{h}{\sigma}} \right) \right]$$

*Fedele F. 2006. Extreme Events in nonlinear random seas. J. of Offshore Mechanics and Arctic Eng., ASME, 128, 11-16.

$$\xi_{nl} = \xi \sqrt[4]{1 + \chi \xi^2}$$

**Symmetric
Third order effects**



$$\tilde{\xi}_{nl} = \xi_{nl} + \frac{\mu}{2} \xi_{nl}^2$$

**Second order
Bound effects**



$$\Pr[\tilde{\xi}_{nl} > y] = \exp\left(-\frac{1}{2} \xi^2(y)\right)$$

$$\tilde{\xi}_{nl} = y$$

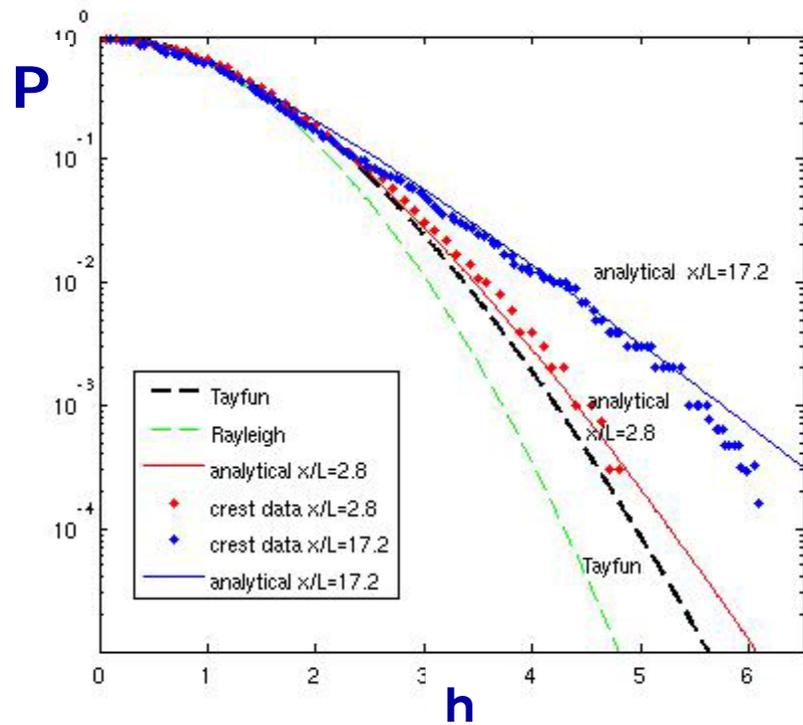
$$\chi = (4 \ln 2) BFI^2 \left(1 - \frac{\cos\left(\frac{1}{2} \arctan \frac{\varepsilon^2 \omega_p T}{BFI^2 2 \ln 2}\right)}{\sqrt[4]{1 + \left(\frac{\varepsilon^2 \omega_p T}{BFI^2 2 \ln 2}\right)^2}} \right)$$

$$BFI = \frac{\sqrt{2} \varepsilon}{\Delta k / k_0}$$

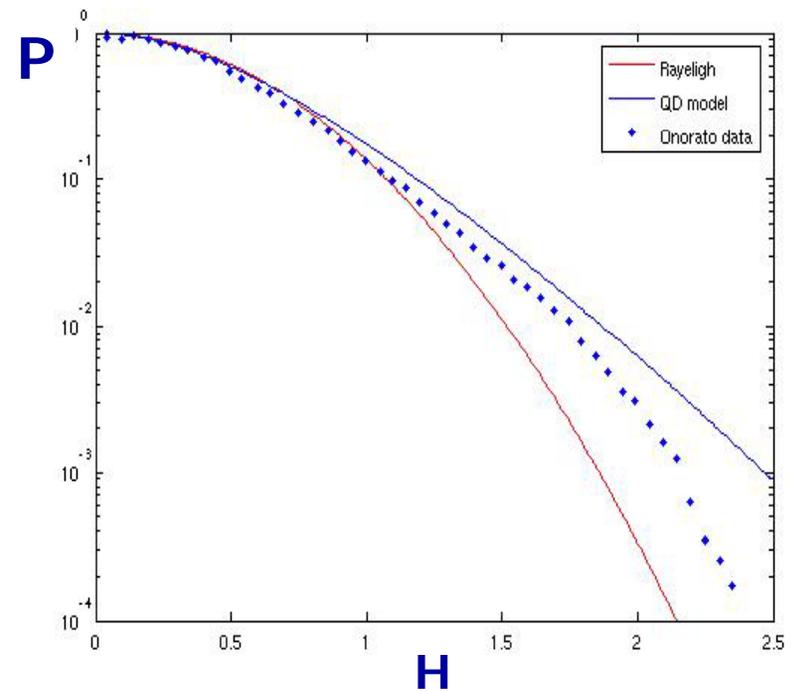
THE PROBABILITY OF EXCEEDING A FIXED WAVE AMPLITUDE

Wave tank experiments:
unidirectional narrow-band seas (Onorato et al. 2005)

THIRD ORDER MODULATION + SECOND ORDER EFFECTS



Wave crest



Wave height

THE PROBABILITY OF EXCEEDANCE cont'd

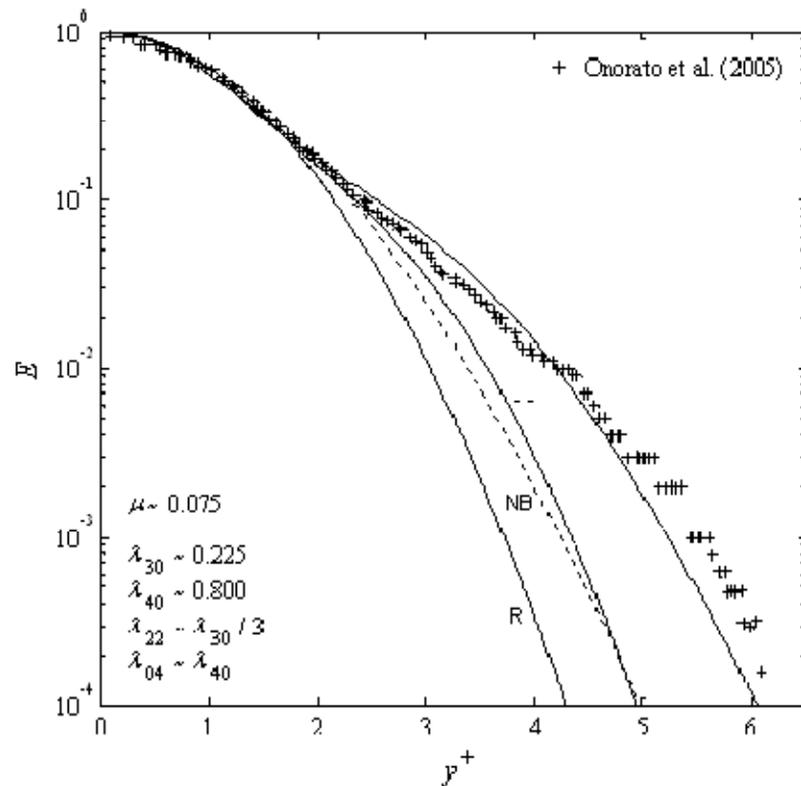
Wave tank experiments:

unidirectional narrow-band seas (Onorato et al. 2005)

TERN platform, time series (6,000 waves)

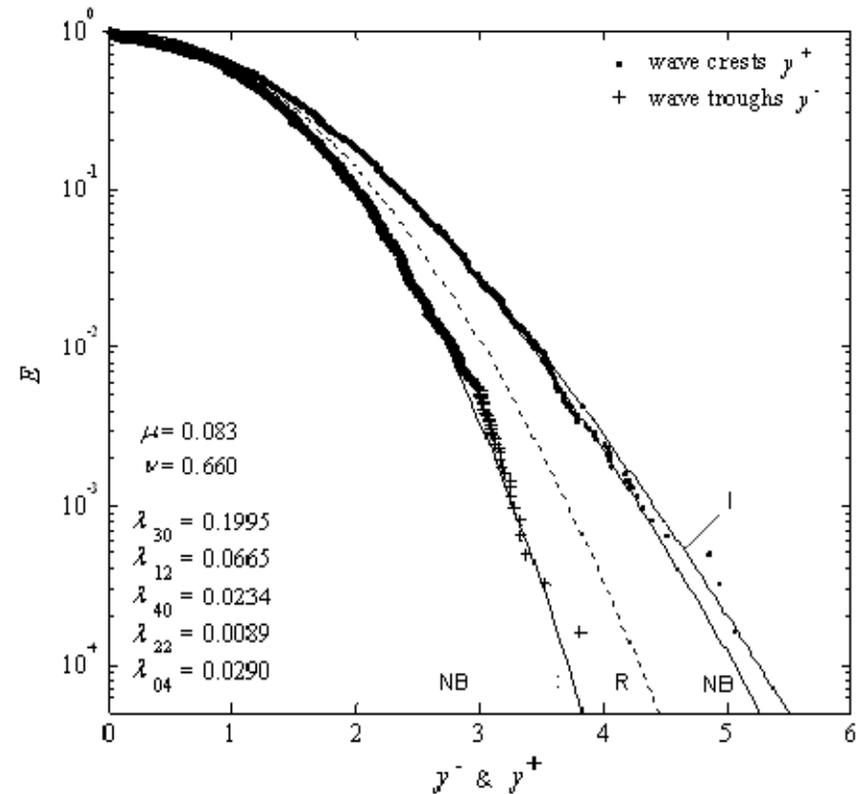
unrealistic ocean conditions

MODULATION + SECOND ORDER EFFECTS



realistic ocean conditions

SECOND ORDER EFFECTS DOMINANT



*Tayfun A. Fedele F., **Wave heights and nonlinear effects**. Ocean Engineering (submitted)

Fedele F., Tayfun A. **Explaining extreme waves by a theory of Stochastic wave groups. PROCEEDINGS of OMAE 2006 (in press)